

Introduction to the stochastic Galerkin approach for the solution of PDEs with parameters or uncertainties

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Recapitulation of Random variable and Probability measure

Structure of the problem, discretization and suitable problems

System of equations and assembly

Approaches to the solution of the system



• 2D Darcy flow problem

$$\begin{cases} -\operatorname{div}\left(k\left(x;\boldsymbol{Z}\right)\cdot\nabla u\left(x;\boldsymbol{Z}\right)\right)=f\left(x;\boldsymbol{Z}\right), & \forall x\in\mathcal{D}\\ u\left(x;\boldsymbol{Z}\right)=u_{D}\left(x;\boldsymbol{Z}\right), & \forall x\in\Gamma_{D},\\ \frac{\partial u(x;\boldsymbol{Z})}{\partial n(x)}=u_{N}\left(x;\boldsymbol{Z}\right), & \forall x\in\Gamma_{N} \end{cases}$$

Z can be understood as $Z : \Omega \to \mathbb{R}^M$ and k(x; Z), $u(x; Z) : \mathcal{D} \times \mathbb{R}^M \to \mathbb{R}$



- **non-intrusive** approaches examines only point values of the solution (much easier to deploy)
 - PDE solution in fixed values of the parameter is handled as **black-box** by an existing software
 - Collocation or Monte Carlo methods
- stochastic Galerkin (SG) method is an intrusive approach
 - needs new framework for the solution of given problems
 - it provides error estimates

Recapitulation of Random variable and Probability measure

Random variables and probability measure



- continuous random variable Z is a map from sample space $\Omega \to \mathbb{R}$
- can be described by probability density $f : \mathbb{R} \to \mathbb{R}^+_0$, $\int f(x) dx = 1$
- distribution of a random variable defines a probability measure
 - in the case of a continuous random variable

$$\int_{\mathbb{R}} \dots dFZ = \int_{\mathbb{R}} \dots f(x) dx = \mathbb{E}_{Z} (\dots)$$



Introduction to the stochastic Galerkin



- is a vector of M (let assume continuous) random variables $\boldsymbol{Z} = (Z_1, \dots, Z_M)$
- can be described by joint probability density $f_Z : \mathbb{R}^M \to \mathbb{R}^+_0$, $\int_{\mathbb{R}^M} f(x) \, dx = 1$
- random vector of independent random variables has joint probability density in form

$$f_{\boldsymbol{Z}}\left(\boldsymbol{z}
ight)=\prod_{i=1}^{M}f_{Z_{i}}\left(z_{i}
ight)$$



• square integrable functions of a random vector \boldsymbol{Z} creates space

$$L^{2}_{\mathsf{dFZ}}\left(\mathbb{R}^{M}\right) := \left\{ f: \mathbb{R}^{M} \to \mathbb{R}: \int_{\mathbb{R}^{M}} f\left(\mathbf{z}\right)^{2} dF\mathbf{Z} < \infty \right\}$$

• with inner product

$$(u,v) := \int_{\mathbb{R}^{M}} u(z) \cdot v(z) \, dFZ = \mathbb{E}_{Z} \left(u(Z) \cdot v(Z) \right)$$



- orthogonal polynomial basis of L^2_{dFZ}
- polynomial chaos with respect to a random variable Z
- set of polynomials of increasing degree ψ_i (Z) (i denotes degree of polynomial)
- orthogonal = satisfying $\mathbb{E}_{Z} (\psi_{i}(Z) \psi_{j}(Z)) = \gamma_{i} \delta_{ij}$, where $\gamma_{i} = \mathbb{E}_{Z} (\psi_{i}(Z)^{2})$
- generates basis of $L^{2}_{\mathrm{d}\mathit{FZ}}\left(\mathbb{R}\right)$

Introduction to the stochastic Galerkin

Polynomial chaos Nd

- assume random vector **Z** of **independent** random variables
- polynomial chaos for $L^2_{dEZ}(\mathbb{R}^M)$ can be created by a product of 1d polynomials of Z_i

$$\Psi_{i}(\boldsymbol{Z})=\prod_{k=1}^{M}\psi_{i_{k}}(Z_{k}),$$

where i denotes the multi-index of size M.

tensor product polynomials of maximal degree *n* are polynomials with

$$i_k \leq n, \forall k = 1, \ldots, M$$

complete polynomials of maximal degree *n* are polynomials with

$$|i| = \sum_{k=1}^{M} i_k \le n$$



Complete vs. tensor polynomials





Figure 1: Error dependence on the max. degree of the PC basis



Figure 2: Error dependence on the size of the PC basis

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Let $\{\psi_i\}, \psi_{-1} = 0, \psi_0 = 1$ be polynomials related to probability measure dFZ (single random variable), then

$$\psi_{i+1}(Z) = (Z - \alpha_i) \psi_i(Z) - \beta_i \psi_{i-1}(Z),$$

$$\alpha_i = \frac{\mathbb{E}_Z (Z\psi_i(Z) \psi_i(Z))}{\mathbb{E}_Z (\psi_i(Z) \psi_i(Z))}, \beta_i = \frac{\mathbb{E}_Z (\psi_i(Z) \psi_i(Z))}{\mathbb{E}_Z (\psi_{i-1}(Z) \psi_{i-1}(Z))}$$

- evaluation of orthogonal polynomials of higher degree can be numerically unstable
 - do not evaluate weights of the polynomials (a₀ + a₁x + ...), numerically unstable (infeasible for degree 20 and more in int32)
 - recurrence formula for computing the point values of the polynomial is favorable

MB

Hermite polynomials

• polynomial chaos basis for standard normal random variable consist of Hermite polynomials

$$\psi_{i+1}(Z) = 2Z\psi_i(Z) - i\psi_{i-1}(Z)$$



Introduction to the stochastic Galerkin



Structure of the problem, discretization and suitable problems



• 2D Darcy flow problem

$$\begin{cases} -\operatorname{div}\left(k\left(x;\boldsymbol{Z}\right)\cdot\nabla u\left(x;\boldsymbol{Z}\right)\right)=f\left(x;\boldsymbol{Z}\right), & \forall x\in\mathcal{D}\\ u\left(x;\boldsymbol{Z}\right)=u_{D}\left(x;\boldsymbol{Z}\right), & \forall x\in\Gamma_{D},\\ \frac{\partial u(x;\boldsymbol{Z})}{\partial n(x)}=u_{N}\left(x;\boldsymbol{Z}\right), & \forall x\in\Gamma_{N} \end{cases}$$

Z can be understood as $Z : \Omega \to \mathbb{R}^M$ and k(x; Z), $u(x; Z) : \mathcal{D} \times \mathbb{R}^M \to \mathbb{R}$



- randomness in f and boundary conditions affects the right hand side
 - if k is deterministic, we obtain decoupled system = matrix of the resulting system of equations is block diagonal
- randomness in material k affects the matrix of the resulting system of equations
 - we obtain coupled system



• 2D Darcy flow problem

$$\begin{cases} -\operatorname{div}\left(k\left(x;\boldsymbol{Z}\right)\cdot\nabla u\left(x;\boldsymbol{Z}\right)\right)=f\left(x\right), & \forall x\in\mathcal{D}\\ u\left(x;\boldsymbol{Z}\right)=u_{D}\left(x\right), & \forall x\in\Gamma_{D},\\ \frac{\partial u\left(x;\boldsymbol{Z}\right)}{\partial n\left(x\right)}=u_{N}\left(x\right), & \forall x\in\Gamma_{N} \end{cases}$$

Z can be understood as $Z : \Omega \to \mathbb{R}^M$ and k(x; Z), $u(x; Z) : \mathcal{D} \times \mathbb{R}^M \to \mathbb{R}$



- we can view k (x; Z), u (x; Z) from the perspective of Monte Carlo sampling
 - sample of Z (e.g. \tilde{Z}) gives us a sample of $k\left(x;\tilde{Z}\right) \in L^{\infty}\left(\mathcal{D}\right)$ and $u\left(x;\tilde{Z}\right) \in H^{1}\left(\mathcal{D}\right)$
- than $u(x; \mathbf{Z}) \in L^2_{\mathsf{d}F\mathbf{Z}}\left(\mathbb{R}^M, H^1(\mathcal{D})\right)$

$$L^{2}_{\mathsf{dFZ}}\left(\mathbb{R}^{M},H^{1}\left(\mathcal{D}\right)\right):=\left\{f:\mathbb{R}^{M}\rightarrow H^{1}\left(\mathcal{D}\right):\int_{\mathbb{R}^{M}}\left\|f\left(\mathbf{Z}\right)\right\|^{2}_{H^{1}\left(\mathcal{D}\right)}\,\mathsf{dFZ}<\infty\right\}$$

• $L^{2}_{dFZ}(\mathbb{R}^{M}, H^{1}(\mathcal{D}))$ is isometrically isomorphic with tensor product space $H^{1}(\mathcal{D}) \otimes L^{2}_{dFZ}(\mathbb{R}^{M})$



- we want to find the solution in the space $V = V_D \otimes V_S$ $(V_D = H^1(\mathcal{D}), V_S = L^2_{dFZ}(\mathbb{R}^M))$
- we need to solve the following

$$\begin{cases} \text{Find } u = \sum_{i=1}^{\infty} u_{D,i} \cdot u_{S,i} \in V_D \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ a(u,v) = b(v), \, v = v_D \cdot v_S \end{cases}, \quad (v,v) \in V_D \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ (v,v) = b(v), \, v = v_D \cdot v_S \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ (v,v) = b(v), \, v = v_D \cdot v_S \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ (v,v) = b(v), \, v = v_D \cdot v_S \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ (v,v) = b(v), \, v = v_D \cdot v_S \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ (v,v) = b(v), \, v = v_D \cdot v_S \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ (v,v) = b(v), \, v = v_D \cdot v_S \otimes V_S, \, \forall v_D \in V_D, v_S \in V_S : \\ (v,v) = b(v), \, v = v_D \cdot v_S \otimes V_S, \, \forall v_D \in V_D, v_S \otimes V_S, \, \forall v_D \in V_D, v_S \otimes V_S, \, \forall v_D \in V_D, v_S \otimes V_S \otimes V_S, \, \forall v_D \in V_D, \, v_S \otimes V_S, \,$$

$$a(u, v) := a(u_D \cdot u_S, v_D \cdot v_S) = \sum_{r=1}^R a_{D,r}(u_D, v_D) \cdot a_{S,r}(u_S, v_S)$$

$$b(v) := b(v_D \cdot v_S) = \sum_{r=1}^{S} b_{D,r}(v_D) \cdot b_{S,r}(v_S)$$



$$a(u, v) := \int_{\mathbb{R}^{M}} \int_{\mathcal{D}} k(x; \mathbf{Z}) \cdot \nabla_{x} u(x; \mathbf{Z}) \cdot \nabla_{x} v(x; \mathbf{Z}) dx dF\mathbf{Z},$$
$$b(v) := \int_{\mathbb{R}^{M}} \int_{\mathcal{D}} f(x) \cdot v(x; \mathbf{Z}) dx dF\mathbf{Z},$$

- we use discretization as tensor product of basis functions on domain and basis functions on sample space $\langle \varphi_1(x), \ldots, \varphi_{N_D}(x) \rangle \otimes \langle \psi_1(\mathbf{Z}), \ldots, \psi_{N_P}(\mathbf{Z}) \rangle \subset V_H :$ $\langle \varphi_1(x), \ldots, \varphi_{N_D}(x) \rangle \subset H_0^1(\mathcal{D})$ and $\langle \psi_1(\mathbf{Z}), \ldots, \psi_{N_P}(\mathbf{Z}) \rangle \subset L^2_{\mathsf{dFZ}}(\mathbb{R}^M)$
- As basis functions on physical domain we can take standard finite elements functions (piece-wise linear)
- For basis functions on sample (random parameter) space we choose polynomial chaos basis



• recall the

$$a(u,v) := \int_{\mathbb{R}^{M}} \int_{\mathcal{D}} k(x; \mathbf{Z}) \cdot \nabla_{x} u(x; \mathbf{Z}) \cdot \nabla_{x} v(x; \mathbf{Z}) \, \mathrm{d}x \, \mathrm{d}F\mathbf{Z}$$

• if we have separable form of the material field

$$k(x; \mathbf{Z}) = \sum_{m=1}^{N} k_m(x) \cdot g_m(\mathbf{Z})$$

• than resulting system of equation is also separable

$$\left(\mathbb{A}\right)_{ij,kl} = \sum_{m=1}^{N} \left(\int_{\mathcal{D}} k_m(x) \nabla \varphi_j(x) \nabla \varphi_l(x) \, \mathrm{d}x \right) \cdot \left(\int_{\mathbb{R}^M} g_m(\mathbf{Z}) \psi_i(\mathbf{Z}) \, \psi_k(\mathbf{Z}) \, \mathrm{d}F\mathbf{Z} \right),$$



• it is natural for the Karhunen-Loève decomposition of the random field

$$k(x; \mathbf{Z}) = \sum_{m=1}^{M} k_i(x) \cdot Z_i$$

or material field with given interfaces

• in other cases it can be obtained by a projection into polynomial chaos

$$k(x; \mathbf{Z}) \approx \sum_{m=1}^{N} k_i(x) \cdot \Psi_i(\mathbf{Z})$$

System of equations and assembly



• After discretization we get following Galerkin system of equations:

$$\mathbb{A}\cdot\overline{u_H}=\overline{b},$$

$$\begin{split} (\mathbb{A})_{ij,kl} &= \sum_{m=1}^{N} \int_{\mathcal{D}} k_{m}(x) \, \nabla \varphi_{j}(x) \, \nabla \varphi_{l}(x) \, \mathrm{d}x \cdot \int_{\mathbb{R}^{M}} g_{m}(\boldsymbol{Z}) \, \psi_{i}(\boldsymbol{Z}) \, \psi_{k}(\boldsymbol{Z}) \, \mathrm{d}F\boldsymbol{Z}, \\ &\mathbb{A} = \sum_{m=1}^{N} G_{m} \otimes K_{m}, \\ &(K_{m})_{ij} = \int_{\mathcal{D}} k_{m}(x) \, \nabla \varphi_{i}(x) \cdot \nabla \varphi_{j}(x) \, \mathrm{d}x, \\ &(G_{m})_{ij} = \int_{\mathbb{R}^{M}} g_{m}(\boldsymbol{Z}) \, \psi_{i}(\boldsymbol{Z}) \, \psi_{j}(\boldsymbol{Z}) \, \mathrm{d}F\boldsymbol{Z}. \end{split}$$



$$\overline{b} = \overline{g} \otimes \overline{f}$$
$$(\overline{g})_{i} = \int_{\mathbb{R}^{M}} \psi_{i}(\mathbf{Z}) \, \mathrm{d}F\mathbf{Z}$$

in this case \overline{g} has only one nonzero value in the first position

$$\left(\overline{f}\right)_{i} = \int_{\mathcal{D}} f(x) \cdot \varphi_{i}(x) \, \mathrm{d}x$$

 Note: with non-homogeneous boundary conditions the r.h.s will have the general form

$$\overline{b} = \sum_{i=1}^{NN} \overline{g_i} \otimes \overline{f_i}$$

Matrix: stiffness matrix of blocks





Figure 3: Permuted system matrix for physical domain discretization 4x4 and 10 polynomial chaos basis function

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Matrix: blocks of stiffness matrix





Figure 4: System matrix for physical domain discretization 4x4 and 10 polynomial chaos basis function

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Examples of sparsity pattern





Figure 5: Sparsity pattern of $G_i, g_i = \exp(Z_i)$

Examples of sparsity pattern





Figure 6: Sparsity pattern of $\sum_i G_i$

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Assembly of G_i for $g_i = \exp(Z_i)$



assume normalized Hermite polynomials $\psi_i(\mathbf{Z}) = \prod_{j=1}^{M} \psi_{i,j}(Z_j)$

$$(G_m)_{k,l} = \int_{\mathbb{R}^M} \psi_k(\mathbf{Z}) \psi_l(\mathbf{Z}) \exp(a_m Z_m + b_m) dF\mathbf{Z} = \int_{\mathbb{R}} \psi_{k,m}(Z_m) \psi_{l,m}(Z_m) \exp(a_m Z_m + b_m) dFZ_m \prod_{i \neq m} \int_{\mathbb{R}} \psi_{k,i}(Z_i) \psi_{l,i}(Z_i) dFZ_i,$$

where $\prod_{i \neq m} \int_{\mathbb{R}} \psi_{k,i}(Z_i) \psi_{l,i}(Z_i) dFZ_i$ is either 0 if some of the polynomials $\psi_{k,i}$ do not equal $\psi_{l,i}$ or 1 otherwise. The rest of the formula can be evaluated according to the following

$$\int_{\mathbb{R}} f(Z) \exp(aZ + b) dFZ = \int_{\mathbb{R}} f(Z) \exp(aZ + b) \frac{1}{\sqrt{2\pi}} \exp\left(-\frac{Z^2}{2}\right) dZ = \exp\left(\frac{a^2}{2} + b\right) \int_{\mathbb{R}} f(Z) \frac{1}{\sqrt{2\pi}} \exp\left(-\frac{(Z-a)^2}{2}\right) dZ = \exp\left(\frac{a^2}{2} + b\right) \int_{\mathbb{R}} f(Z) dF\tilde{Z},$$

where $\tilde{Z} \sim \mathcal{N}(a, 1)$.

 $\int_{\mathbb{R}} f(Z) dF\tilde{Z}$ can be precisely computed using the Gauss-Hermite quadrature rule or evaluated analytically (if f is polynomial), but the analytic formula can be numerically unstable.



$$\mathbb{A}=\sum_{m=1}^N G_m\otimes K_m,$$

Matrix multiplication

$$\mathbb{A}x = \operatorname{vec}\left(\sum_{m=1}^{N} K_m \tilde{x} G_m^{\mathsf{T}}\right),$$

where \tilde{x} is reshaped x into a matrix of corresponding dimensions.

Approaches to the solution of the system



- Krylov method with a suitable preconditioner with a low-rank compression in each iteration [7, 2, 4, 6]
 - preconditioners: block diagonal [15], Kronecker [18] hierarchical Schur [17]
- Generalized Spectral Decomposition, where the problem is treated as an extended eigenvalue problem and the solution is built using the power-type method [10, 9, 14, 11, 13, 12] (even nonlinear problems)
- low rank tensor approximation (also for multiple independent parameters): approximates the solution using the tensor train or hierarchical Tucker format [5, 1, 3].
- reduced basis (RB) approach construction of the RB using the rational Krylov approximations [16]; using MC sampling [19], using sampling via sparse grids is presented in [8]



• assume a low rank approximation of the solution

$$\boldsymbol{x} \approx \boldsymbol{x}_{\boldsymbol{k}} = W_k \boldsymbol{y}_{\boldsymbol{k}},$$

where $W_k = [w_1, \ldots, w_k] \in \mathbb{R}^{N_d \times k}$ is a **Reduced basis** of size k (k orthonormal vectors)

- here W_k approximates "deterministic" part of the solution x (given by matrices K_m)
- matrix y_k can be obtained using the Galerkin condition on the residual of x_k

$$R_k := \sum_{m=1}^M K_m(W_k y_k) G_m^T - \sum_{m=1}^M f_m g_m^T$$

$$W_k^T R_k = 0 \Rightarrow \sum_{m=1}^M W_k^T K_m W_k y_k G_m^T = \sum_{m=1}^M W_k^T f_m g_m^T$$

which is a system of equations of much lower dimension (can be easily solved by e.g. PCG)

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Reduced basis solver

Low rank approximation of the solution



• assume a low rank approximation of the solution

$$\mathbf{x} \approx \mathbf{x}_{\mathbf{k}} = W_k \mathbf{y}_{\mathbf{k}},$$

where $W_k = [w_1, \ldots, w_k] \in \mathbb{R}^{N_d \times k}$ is a **Reduced basis** of size k (k orthonormal vectors)

- here W_k approximates "deterministic" part of the solution x (given by matrices K_m)
- matrix y_k can be obtained using the Galerkin condition on the residual of x_k

$$R_{k} := \sum_{m=1}^{M} K_{m} (W_{k} y_{k}) G_{m}^{T} - \sum_{m=1}^{M} f_{m} g_{m}^{T}$$
$$W_{k}^{T} R_{k} = 0 \Rightarrow \sum_{m=1}^{M} W_{k}^{T} K_{m} W_{k} y_{k} G_{m}^{T} = \sum_{m=1}^{M} W_{k}^{T} f_{m} g_{m}^{T}$$

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Reduced rational Krylov subspaces



- series of symmetric positive definite (SPD) matrices $\{K_m\}_{m=1,...,M}$ and a nonzero vector v
- the first iteration generates addition to the basis:

$$\left\langle K_{1}^{-1}v,\ldots,K_{M}^{-1}v\right\rangle$$

• the second iteration generate addition to the basis:

$$\langle K_1^{-1}K_1^{-1}v, K_1^{-1}K_2^{-1}v, \dots, K_M^{-1}K_{M-1}^{-1}v, K_M^{-1}K_M^{-1}v \rangle$$

Building the reduced basis



- starting vector from $\hat{f_m} = K_0^{-1} f_m$
- using matrices $\{K_0^{-1}(K_m + \alpha_m K_0)\}$
- full RRKS scheme is impractical because in each subsequent iteration we need to construct larger bases
- the remedy to this is to iteratively select a vector v from the current basis and expand the basis only by the first iteration of the RRKS
- during the orthogonalisation step, calculate weights corresponding to the norm of independent part of the vector
 - next direction according to the calculated weights
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Figure 1: different values of α (degree 4, grid level 10)

RB convergence - polynomial degree





Figure 2: different maximal degree of polynomial chaos (grid lvl 10)





Figure 3: different grid levels (polynomial degree 4)

Monte Carlo sampling

We compare two random sampling procedures

- Crude MC
- Sampling using random process, which avoids already generated samples (MCMC)
 - assume sampling according to unscaled probability density f (x)
 - each consecutive sample X_i is generated according to unscaled probability density

$$\bar{f}(x) = f(x) \cdot \min\left\{1 - \exp\left(-\frac{1}{2}\left(x - X_j\right)^T \cdot \Sigma^{-1} \cdot \left(x - X_j\right)\right)\right\}_{j=1\dots i-j}$$

 these samples are generated using Metropolis Hastings method (we can take random sample of the chain, or the one with the highest value of f
 (x))

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Unscaled probability density after 10 samples





Figure 4: Unscaled probability density after 10 samples ($\Sigma = 10 \cdot I$)





Figure 5: left MCMC, right MC

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Loop until satisfactory reduced basis is obtained:

- 1. Propose *n* samples (MC or MCMC)
- 2. Calculate residual of reduced solution in this samples
- 3. Solve systems, using best N samples (largest residual error)
- Add these N samples into the reduced basis (truncate if needed)





Figure 6: different number of proposal samples n

MC vs. MCMC number of parallel solutions





Figure 7: different number of parallel solutions N





Figure 8: different maximal degree of polynomial chaos (grid lvl 10)





Figure 9: different grid levels (polynomial degree 4)





Figure 10: grid lvl 10, polynomial chaos degree 4



MCMC:

- better convergence (same rate, but shifted)
- easier to implement (it is possible to call black box solvers)
- random (different runs can differ)

RRKS:

- it solves the same systems (much cheaper assembly of systems and preconditioners)
- depends on the choice of α (can be difficult)

Deflated conjugate gradients



- DCG method is an extension of the standard conjugate gradient (CG or PCG if using a preconditioner) method
- DCG method takes an additional parameter in the form of the deflation basis *W* (linearly independent, we will use orthonormal).
- the deflation basis *W* should be able to describe the sought solution reasonably well
- the DCG method looks for the solution outside of the deflation basis *W*

- DCG diference in comparison to the conjugate gradient implementation:
 - the residual of the initial guess should be orthogonal on W (e.g. by using $x_0 = W (W^T A W)^{-1} W^T b$)
 - in each iteration the DCG projects the preconditioned residual $\tilde{z}_i = P z_i$ using the projector $P = I W (W^T A W)^{-1} W^T A$
- in our case the reduced basis W_i in *i*-th iteration of RB solver can be directly used as the deflation basis
- only significant additional cost compared to the standard CG method is the solution of systems with $Q_i = W_i^T A W_i$ (in *i*-th iteration of the RB solver)
 - In our application, the size of the matrix Q_i (corresponding to a size of the RB) is reasonably small and we can use e.q. an explicit inversion.





- we test the solution of systems $(K_m + \alpha K_0)$ using: Schwarz, diagonal and ichol preconditioner
 - the Schwarz preconditioner was set using 30 subdomains = equal column slices of our square domain of the size 1/20
 - the incomplete Cholesky preconditioner was build with no filling allowed

RRKS: DCG - results (number of iterations CG/DCG)





Figure 11: Num. of the DCG iterations compared to the PCG (mean DCG/PCG iterations per RB iter.)

MCMC: DCG - results (number of iterations CG/DCG)





Figure 12: Num. of the DCG iterations compared to the PCG (mean DCG/PCG iterations per RB iter.)



	Ad. Schwarz	diagonal	ichol (nofill)
RRKS savings in %	72.32%	73.47%	73.33%
MCMC savings in %	82.58%	83.48%	83.06%

 Table 1: Computational savings using DCG with the RB as a deflation basis