

Exponential growth, exponential scaling, and the discontinuous Galerkin method

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- 1 Gronwall's lemma
- 2 Exponential scaling
- 3 Discontinuous Galerkin
- 4 Error analysis
- 5 Construction of μ

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Gronwall's lemma

If for all $t \in (0, T)$

$$u'(t) \leq \beta(t)u(t)$$

then

$$u(t) \leq u(0)e^{\int_0^t \beta(s)ds}$$

Gronwall's lemma - integral form

Let $\beta \geq 0$. If for all $t \in (0, T)$

$$u(t) \leq \alpha + \int_0^t \beta(s)u(s)ds$$

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$$u(t) \leq \alpha e^{\int_0^t \beta(s)ds}$$

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Abstract Gronwall's lemma

Let M be a sequentially closed topological space which has a partial ordering \leq . Let $A : M \rightarrow M$ such that

- $x \leq y \Rightarrow A(x) \leq A(y)$.
- A has a unique fixed point x^* .
- $A^n(x) \rightarrow x^*$ for all x .

Then

$$x \leq A(x) \Rightarrow x \leq x^*.$$

Proof:

- $x \leq A(x)$
- $x \leq A^n(x)$
- Limit $\Rightarrow x \leq x^*$

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Advection-reaction

$$\frac{\partial u}{\partial t} + \mathbf{a} \cdot \nabla u + cu = 0.$$

Estimates of u : multiply by u , \int_{Ω}

$$\int_{\Omega} \frac{\partial u}{\partial t} u \, dx + \int_{\Omega} \mathbf{a} \cdot \nabla u u \, dx + \int_{\Omega} cu^2 \, dx = 0.$$

By Green's theorem

$$\int_{\Omega} \mathbf{a} \cdot \nabla u u \, dx = \frac{1}{2} \int_{\partial\Omega} \mathbf{a} \cdot \mathbf{n} u^2 \, d\sigma - \frac{1}{2} \int_{\Omega} \operatorname{div} \mathbf{a} u^2 \, dx.$$

Therefore

$$\frac{1}{2} \frac{d}{dt} \|u\|^2 + \frac{1}{2} \int_{\partial\Omega} \mathbf{a} \cdot \mathbf{n} u^2 \, d\sigma + \int_{\Omega} \left(c - \frac{1}{2} \operatorname{div} \mathbf{a} \right) u^2 \, dx = 0.$$

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Continuous problem

$$\frac{1}{2} \frac{d}{dt} \|u\|^2 + \int_{\Omega} (c - \frac{1}{2} \operatorname{div} a) u^2 dx \leq 0.$$

- Elliptic case: $c - \frac{1}{2} \operatorname{div} a \geq 0$

$$\frac{d}{dt} \|u\|^2 \leq 0.$$

- Non-elliptic case: $c - \frac{1}{2} \operatorname{div} a \leq c_0$

$$\frac{1}{2} \frac{d}{dt} \|u\|^2 \leq c_0 \|u\|^2.$$

Gronwall gives

$$\|u(t)\|^2 \leq e^{2c_0 t} \|u(0)\|^2.$$

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Exponential scaling I.

$$\frac{\partial u}{\partial t} + \mathbf{a} \cdot \nabla u + cu = 0.$$

Exponential scaling:

$$u(x, t) = e^{\alpha t} w(x, t).$$

This gives

New ellipticity condition:

$$\alpha + c - \frac{1}{2} \operatorname{div} \mathbf{a} \geq 0$$

Useless in our case!

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Exponential scaling II.

$$\frac{\partial u}{\partial t} + \mathbf{a} \cdot \nabla u + cu = 0.$$

Exponential scaling (Nävert, 1982): Choose $\mu_0 \in \mathbb{R}^d$,

$$u(x, t) = e^{\mu_0 \cdot x} w(x, t).$$

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Vector field \mathbf{a} must point in one direction!

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Exponential scaling III.

$$\frac{\partial u}{\partial t} + \mathbf{a} \cdot \nabla u + cu = 0.$$

Exponential scaling (Ayuso, Marini, 2009): Choose $\mu : \Omega \rightarrow \mathbb{R}$,

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Lemma – Devinatz, Ellis, Friedman (1974)

$\exists \mu : \mathbf{a} \cdot \nabla \mu \geq \gamma_0 > 0 \iff \mathbf{a}$ possesses neither closed curves nor stationary points.

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Exponential scaling IV.

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Exponential scaling (Kučera, Shu): Choose $\mu : \Omega \times (0, T) \rightarrow \mathbb{R}$,

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Weak form of exponential scaling

$$\int_{\Omega} \frac{\partial u}{\partial t} v + \mathbf{a} \cdot \nabla u v + c u v \, dx = 0.$$

- Set $u(x, t) = e^{\mu(x, t)} \tilde{u}(x, t)$.
- Set $\hat{v}(x, t) = e^{-\mu(x, t)} v(x, t)$:

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Continuous problem

We seek $u : Q_T = \Omega \times (0, T) \rightarrow \mathbb{R}$ such that

$$\begin{aligned} \frac{\partial u}{\partial t} + \mathbf{a} \cdot \nabla u + cu &= 0 && \text{in } Q_T, \\ u &= u_D && \text{on } \partial\Omega^- \times (0, T), \\ u(x, 0) &= u^0(x), && x \in \Omega. \end{aligned}$$

- \mathcal{T}_h is a triangulation of Ω .
- $S_h = \{v_h; v_h|_K \in P^p(K), \forall K \in \mathcal{T}_h\}$,

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DG discretization

We seek $u_h \in C^1(0, T; S_h)$ such that

$$\left(\frac{\partial u_h}{\partial t}, v_h \right) + b_h(u_h, v_h) + c_h(u_h, v_h) = I_h(v_h), \quad \forall v_h \in S_h,$$

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We seek $u_h \in C^1(0, T; S_h)$ such that

$$\left(\frac{\partial u_h}{\partial t}, v_h \right) + b_h(u_h, v_h) + \mathbf{c}_h(u_h, v_h) = I_h(v_h), \quad \forall v_h \in S_h,$$

Reaction form

$$c_h(u, v) = \int_{\Omega} cuv \, dx$$

$$\frac{\partial u}{\partial t} + \mathbf{a} \cdot \nabla u + cu = 0.$$

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Advection form

$$\begin{aligned} b_h(u, v) = & \sum_{K \in \mathcal{T}_h} \int_K (\mathbf{a} \cdot \nabla u) v \, dx - \sum_{K \in \mathcal{T}_h} \int_{\partial K^- \setminus \partial \Omega} (\mathbf{a} \cdot \mathbf{n}) [u] v \, dS \\ & - \sum_{K \in \mathcal{T}_h} \int_{\partial K^- \cap \partial \Omega} (\mathbf{a} \cdot \mathbf{n}) u v \, dS, \end{aligned}$$

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Right-hand side form

$$I_h(v) = - \sum_{K \in \mathcal{T}_h} \int_{\partial K^- \cap \partial \Omega} (\mathbf{a} \cdot \mathbf{n}) u_D v \, dx.$$

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- Error $e_h(t) := u(t) - u_h(t) = \eta_h(t_n) + \xi_h(t)$.
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Assumptions on μ

$$0 \leq \mu(x, t) \leq \mu_{\max},$$

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Lemma

Let $\phi = e^{-\mu} \tilde{\xi}$, then

$$\|\Pi_h \phi(t) - \phi(t)\|_{L^2(K)} \leq Ch \|\tilde{\xi}(t)\|_{L^2(K)},$$

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$$\begin{aligned} \left(\frac{\partial \xi}{\partial t}, \Pi_h \phi \right) &= \left(\frac{\partial \xi}{\partial t}, \phi \right) = \left(e^\mu \frac{\partial \tilde{\xi}}{\partial t} + e^\mu \frac{\partial \mu}{\partial t} \tilde{\xi}, e^{-\mu} \tilde{\xi} \right) \\ &= \frac{1}{2} \frac{d}{dt} \|\tilde{\xi}\|^2 + \left(\frac{\partial \mu}{\partial t} \tilde{\xi}, \tilde{\xi} \right). \end{aligned}$$

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Lemma: Let $\frac{\partial \mu}{\partial t} + a \cdot \nabla \mu + c - \frac{1}{2} \operatorname{div} a \geq \gamma_0 > 0$, then

$$\begin{aligned} & \left(\frac{\partial \xi}{\partial t}, \Pi_h \phi \right) + b_h(\xi, \phi) + c_h(\xi, \phi) \\ & \geq \frac{1}{2} \frac{d}{dt} \|\tilde{\xi}\|^2 + \gamma_0 \|\tilde{\xi}\|^2 + \frac{1}{2} \sum_{K \in \mathcal{T}_h} \left(\|[\tilde{\xi}]^2\|_{a, \partial K^- \setminus \partial \Omega}^2 + \|\tilde{\xi}\|_{a, \partial K \cap \partial \Omega}^2 \right). \end{aligned}$$

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$$\int_{\Omega} c \xi (\Pi_h \phi - \phi) dx \leq C \|\tilde{\xi}\| h \|\tilde{\xi}\| = Ch \|\tilde{\xi}\|^2.$$

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Altogether

$$\frac{d}{dt} \|\tilde{\xi}(t)\|^2 + 2\gamma_0 \|\tilde{\xi}(t)\|^2 \leq Ch \|\tilde{\xi}(t)\|^2 + Ch^{2p+1} (|u(t)|_{H^{p+1}}^2 + |u_t(t)|_{H^{p+1}}^2).$$

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$$\begin{aligned} \|\tilde{\xi}(t)\|^2 + \gamma_0 \int_0^t \|\tilde{\xi}(\vartheta)\|^2 d\vartheta \\ \leq Ch^{2p+1} (|u^0|_{H^{p+1}}^2 + |u|_{L^2(0,t;H^{p+1})}^2 + |u_t|_{L^2(0,t;H^{p+1})}^2). \end{aligned}$$

$$\|\xi(t)\| = \|e^{\mu(t)} \tilde{\xi}(t)\| \leq e^{\mu_{\max}} \|\tilde{\xi}(t)\|.$$

Theorem

Let $0 \leq \mu \leq \mu_{\max}$, $\mu(t) \in W^{1,\infty}(\Omega)$ and $\mu_t + a \cdot \nabla \mu + c - \frac{1}{2} \operatorname{div} a \geq \gamma_0$ on Q_T , where $\gamma_0 > 0$. Then

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$$\mu_t + \mathbf{a} \cdot \nabla \mu + c - \frac{1}{2} \operatorname{div} \mathbf{a} \geq \gamma_0 > 0.$$

- Find μ_1 such that

$$\frac{\partial \mu_1}{\partial t} + \mathbf{a} \cdot \nabla \mu_1 = 1.$$

- Pathlines:** $S(\cdot; x_0, t_0)$ is the trajectory of a massless particle

$$S(t_0; x_0, t_0) = x_0 \in \overline{\Omega}, \quad \frac{dS(t; x_0, t_0)}{dt} = \mathbf{a}(S(t; x_0, t_0), t).$$

- Then

$$\frac{d\mu_1(S(t; x_0, t_0), t)}{dt} = \left(\frac{\partial \mu_1}{\partial t} + \mathbf{a} \cdot \nabla \mu_1 \right)(S(t; x_0, t_0), t) = 1,$$

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$$\frac{\partial \mu_1}{\partial t} + \mathbf{a} \cdot \nabla \mu_1 = 1.$$

- Pathlines:** $S(\cdot; x_0, t_0)$ is the trajectory of a massless particle

$$S(t_0; x_0, t_0) = x_0 \in \overline{\Omega}, \quad \frac{dS(t; x_0, t_0)}{dt} = \mathbf{a}(S(t; x_0, t_0), t).$$

- Then

$$\frac{d\mu_1(S(t; x_0, t_0), t)}{dt} = \left(\frac{\partial \mu_1}{\partial t} + \mathbf{a} \cdot \nabla \mu_1 \right)(S(t; x_0, t_0), t) = 1,$$

therefore

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- Assumption: Let the time spent in Ω by any particle carried by the flow be bounded by \widehat{T} . Then $0 \leq \mu \leq \mu_{\max}$.
- If we choose

$$\mu(x, t) = \mu_1(x, t) \left(\left| \inf_{Q_T} \left(c - \frac{1}{2} \operatorname{div} a \right)^- \right| + \gamma_0 \right),$$

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Lemma

Let $a \in L^\infty(Q_T)$ be continuous w.r.t. time and Lipschitz continuous w.r.t. space. Let there exist $a_{\min} > 0$ such that

$$-a(x, t) \cdot \mathbf{n} \geq a_{\min}$$

for all $x \in \partial\Omega^-, t \in [0, T]$. Let the life-time of particles carried by the flow be uniformly bounded by \hat{T} . Then μ is Lipschitz continuous and uniformly bounded.

Resulting estimate

$$\|e_h\|_{L^\infty(L^2)} + \sqrt{\gamma_0} \|e_h\|_{L^2(L^2)} \leq C e^{c\hat{T}} h^{p+1/2}.$$

Standard Gronwall estimates

$$\|e_h\|_{L^\infty(L^2)} \leq C e^{cT} h^{p+1/2}.$$

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Conclusions:

- Gronwall's inequality gives exponential growth w.r.t. time.
- General form of exponential scaling.
- Error estimates for DG exponentially growing not in time but in the time particles carried by the flow spend in Ω .
- Effectively, we apply Gronwall's inequality not in the Eulerian framework, but in the Lagrangian (along individual pathlines).
- Nonlinear case possible (linear case + Zhang-Shu technique).

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Thank you for your attention